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# 一类洋流运动方程的显示行波解

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**摘要:**考察一类描述洋流运动的偏微分方程模型, 借助齐次平衡法的思想, 对模型解的形式进行假设。利用改进的Tanh函数法, 将该偏微分方程组约化为相应的非线性代数方程组, 并借助Maple的符号运算功能获得模型的三角函数形式的周期行波解和双曲函数形式的行波解的精确表达式。

**关 键 词:**洋流运动方程; 非线性偏微分方程组; 周期波解; 双曲函数解

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考虑一类重要的洋流运动偏微分方程模型<sup>[1]</sup>

$$\begin{cases} u_t + uu_x + vu_y - v = -h_x \\ v_t + uv_x + vv_y + u = -h_y \\ h_t + (uh)_x + (vh)_y = 0 \end{cases} \quad (1)$$

其中, 函数  $h(x, y, t)$  代表洋流深度, 函数  $u(x, y, t)$ ,  $v(x, y, t)$  分别表示洋流沿  $x, y$  方向的运动速度。方程组(1)作为一个重要的水波方程, 无论从理论角度还是从应用角度, 其行波解的研究具有非常重要意义。首先, M. Renardy<sup>[2]</sup>重点关注其行波解的不稳定行为, 并证明这种不稳定性源自方程的 Hopf 分岔和带有  $O_2$  对称性的 Hopf 分岔。随后, 刘倩等<sup>[3]</sup>利用动力系统的分岔方法讨论方程(1)各种类型有界行波解的存在性条件, 并发现模型的周期光滑行波会随系统参数的变化而逐渐失去光滑性, 最终演变成周期尖波的动力行为。这些有趣的结果有助于从定性分析的角度理解行波解的存在性、稳定性及其动力学演化行为。但是, 从计算角度, 希望能直接给出模型的显示行波解。因此利用改进的Tanh函数法<sup>[4-9]</sup>求解方程, 得到方程(1)的三角函数周期行波解和双曲函数类型行波解的精确表达式。

对方程(1)做行波变换  $u(x, y, t) = U(\xi)$ ,  $v(x, y, t) = V(\xi)$ ,  $h(x, y, t) = H(\xi)$ ,  $\xi = x + ay - ct$ , 得到

$$\begin{cases} -cU' + UU' + aVU' - V = -H' \\ -cV' + UV' + aVV' + U = -aH' \\ -cH' + (UH)' + a(VH)' = 0 \end{cases} \quad (2)$$

其中'表示  $\frac{d}{d\xi}$ , 对式(2)的第三式关于  $\xi$  积分一次,

得

$$-cH + UH + aVH = g \quad (3)$$

其中  $g$  是积分常数, 则

$$H = \frac{g}{aV + U - c} \quad (4)$$

把式(4)代入式(2), 整理可得

$$\begin{cases} -cU' + UU' + aVU' - V = \frac{g(aV' + U')}{(aV + U - c)^2} \\ -cV' + UV' + aVV' + U = \frac{ag(aV' + U')}{(aV + U - c)^2} \end{cases} \quad (5)$$

根据齐次平衡法<sup>[10]</sup>的思想, 并利用改进的Tanh函数展开法, 可假设方程(5)有如下形式的解

$$U(\xi) = \sum_{i=0}^{i=2} a_i \varphi^i(\xi) \quad V(\xi) = \sum_{i=0}^{i=2} b_i \varphi^i(\xi) \quad (6)$$

其中  $a_i, b_i$  为待定常数,  $\varphi(\xi)$  满足 Riccati 方程

$$\varphi' = p + \varphi^2 \quad (7)$$

$p$  是常数, 将方程(6)代入方程(5), 并且利用方程(7), 令方程中所有  $f^i$  的系数为 0, 得到关于  $a_i, b_i, a, c, p$  的非线性代数方程组

$$(ab_2 + a_2)^2(2ab_2^2 + 2a_2b_2) = 0 \quad (8)$$

$$(ab_2 + a_2)^2(2aa_2b_2 + 2a_2^2) = 0 \quad (9)$$

$$(ab_0 + a_0 - c)^2(aa_1b_0p + a_0a_1p - a_1cp - b_0) - g(ab_1p + a_1p) = 0 \quad (10)$$

$$(ab_0 + a_0 - c)^2(ab_0b_1p + a_0b_1p - b_1cp + a_0) - ag(ab_1p + a_1p) = 0 \quad (11)$$

$$2(ab_1 + a_1)(ab_2 + a_2)(2ab_2^2 + 2a_2b_2) + (ab_2 + a_2)^2(3ab_1b_2 + 2a_1b_2 + a_2b_1) = 0 \quad (12)$$

$$2(ab_1 + a_1)(ab_2 + a_2)(2aa_2b_2 + 2a_2^2) + (ab_2 + a_2)^2(aa_1b_2 + 2aa_2b_1 + 3a_1a_2) = 0 \quad (13)$$

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$$\begin{aligned} & (2(ab_0+a_0-c)(ab_2+a_2)+(ab_1+a_1)^2)(2ab_2^2+2a_2b_2)+ \\ & 2(ab_1+a_1)(ab_2+a_2)(3ab_1b_2+2a_1b_2+a_2b_1)+(ab_2+a_2)^2 \\ & (2ab_2^2p+2ab_0b_2+ab_1^2+2a_2b_2p+2a_0b_2+a_1b_1-2b_2c)=0 \end{aligned} \quad (14)$$

$$\begin{aligned} & (2(ab_0+a_0-c)(ab_2+a_2)+(ab_1+a_1)^2)(2aa_2b_2+2a_2^2)+ \\ & 2(ab_1+a_1)(ab_2+a_2)(aa_1b_2+2aa_2b_1+3a_1a_2)+(ab_2+a_2)^2(2aa_2b_2p+aa_1b_1+2aa_2b_0+2a_2^2p+2a_0a_2+a_1^2-2a_2c)=0 \end{aligned} \quad (15)$$

$$\begin{aligned} & (ab_0+a_0-c)^2(aa_1b_1p+2aa_2b_0p+2a_0a_2p+a_1^2p-2a_2cp-b_1)+2(ab_0+a_0-c)(ab_1+a_1)(aa_1b_0p+a_0a_1p-a_1cp-b_0)- \\ & g(2ab_2p+2a_2p)=0 \end{aligned} \quad (16)$$

$$\begin{aligned} & (ab_0+a_0-c)(2ab_0b_2p+ab_1^2p+2a_0b_2p+a_1b_1p-2b_2cp+a_1)+ \\ & +2(ab_0+a_0-c)(ab_1+a_1)(ab_0b_1p+a_0b_1p-b_1cp+a_0)- \\ & -ag(2ab_2p+2a_2p)=0 \end{aligned} \quad (17)$$

$$\begin{aligned} & (ab_0+a_0-c)^2(3ab_1b_2p+ab_0b_1+2a_1b_2p+a_2b_1p+a_0b_1-b_1c+a_2)+2(ab_0+a_0-c)(ab_1+a_1)(2ab_0b_2p+ab_1^2p+2a_0b_2p+a_1b_1p-2b_2cp+a_1)+(2(ab_0+a_0-c)(ab_2+a_2)+(ab_1+a_1)^2)(ab_0b_1p+a_0b_1cp+a_0)-ag(ab_1+a_1)=0 \end{aligned} \quad (18)$$

$$\begin{aligned} & (ab_0+a_0-c)^2(aa_1b_2p+2aa_2b_1p+aa_1b_0+3a_1a_2p+a_0a_1-a_1c-b_2)+2(ab_0+a_0-c)(ab_1+a_1)(aa_1b_1p+2aa_2b_0p+2a_0a_2p+a_1^2p-2a_2cp-b_1)+(2(ab_0+a_0-c)(ab_2+a_2)+(ab_1+a_1)^2)(aa_1b_0p+a_0a_1p-a_1cp-b_0)-g(ab_1+a_1)=0 \end{aligned} \quad (19)$$

$$\begin{aligned} & 2(ab_0+a_0-c)(ab_1+a_1)(2ab_2^2+2a_2b_2)+(2(ab_0+a_0-c)(ab_2+a_2)+(ab_1+a_1)^2)(3ab_1b_2+2a_1b_2+a_2b_1)+2(ab_1+a_1)(ab_2+a_2)(2ab_2^2p+2ab_0b_2+ab_1^2+2a_2b_2p+2a_0b_2+a_1b_1-2b_2c)+(ab_2+a_2)^2(3ab_1b_2p+ab_0b_1+2a_1b_2p+a_2b_1p+a_0b_1-b_1c+a_2)=0 \end{aligned} \quad (20)$$

$$\begin{aligned} & 2(ab_0+a_0-c)(ab_1+a_1)(2aa_2b_2+2a_2^2)+(2(ab_0+a_0-c)(ab_2+a_2)+(ab_1+a_1)^2)(aa_1b_2+2aa_2b_1+3a_1a_2)+2(ab_1+a_1)(ab_2+a_2)(2aa_2b_2p+aa_1b_1+2aa_2b_0+2a_2^2p+2a_0a_2+a_1^2-2a_2c)+(ab_2+a_2)^2(aa_1b_2p+2aa_2b_1p+aa_1b_0+3a_1a_2p+a_0a_1-a_1c-b_2)=0 \end{aligned} \quad (21)$$

$$\begin{aligned} & (ab_0+a_0-c)^2(2ab_2^2+2a_2b_2)+2(ab_0+a_0-c)(ab_1+a_1)(3ab_1b_2+2a_1b_2+a_2b_1)+(2(ab_0+a_0-c)(ab_2+a_2)+(ab_1+a_1)^2)(2ab_2^2p+2ab_0b_2+ab_1^2+2a_2b_2p+2a_0b_2+a_1b_1-2b_2c)+2(ab_1+a_1)(ab_2+a_2)(3ab_1b_2p+ab_0b_1+2a_1b_2p+a_2b_1p+a_0b_1-b_1c+a_2)+(ab_2+a_2)^2(2ab_0b_2p+ab_1^2p+2a_0b_2p+a_1b_1p-2b_2cp+a_1)=0 \end{aligned} \quad (22)$$

$$\begin{aligned} & (ab_0+a_0-c)^2(2aa_2b_2+2a_2^2)+2(ab_0+a_0-c)(ab_1+a_1)(aa_1b_2+2aa_2b_1+3a_1a_2)+(2(ab_0+a_0-c)(ab_2+a_2)+(ab_1+a_1)^2)(2aa_2b_2p+aa_1b_1+2aa_2b_0+2a_2^2p+2a_0a_2+a_1^2-2a_2c)+2(ab_1+a_1)(ab_2+a_2)(aa_1b_2p+2aa_2b_1p+aa_1b_0+3a_1a_2p+a_0a_1-a_1c-b_2)+(ab_2+a_2)^2(aa_1b_1p+2aa_2b_0p+ \end{aligned}$$

$$2a_0a_2p+a_1^2p-2a_2cp-b_1)=0 \quad (23)$$

$$\begin{aligned} & (ab_0+a_0-c)^2(3ab_1b_2+2a_1b_2+a_2b_1)+2(ab_0+a_0-c)(ab_1+a_1)(2ab_2^2p+2ab_0b_2+ab_1^2+2a_2b_2p+2a_0b_2+a_1b_1-2b_2c)+(2(ab_0+a_0-c)(ab_2+a_2)+(ab_1+a_1)^2)(3ab_1b_2p+ab_0b_1+2a_1b_2p+a_2b_1p+a_0b_1-b_1c+a_2)+2(ab_1+a_1)(ab_2+a_2)(2ab_0b_2p+ab_1^2p+2a_0b_2p+a_1b_1p-2b_2cp+a_1)+(ab_2+a_2)^2(ab_0b_1p+a_0b_1p-b_1cp+a_0)=0 \end{aligned} \quad (24)$$

$$\begin{aligned} & (ab_0+a_0-c)^2(aa_1b_2+2aa_2b_1+3a_1a_2)+2(ab_0+a_0-c)(ab_1+a_1)(2aa_2b_2p+aa_1b_1+2aa_2b_0+2a_2^2p+2a_0a_2+a_1^2-2a_2c)+(2(ab_0+a_0-c)(ab_2+a_2)+(ab_1+a_1)^2)(aa_1b_2p+2aa_2b_1p+aa_1b_0+3a_1a_2p+a_0a_1-a_1c-b_2)+2(ab_1+a_1)(ab_2+a_2)(aa_1b_1p+2aa_2b_0p+2a_0a_2p+a_1^2p-2a_2cp-b_1)+(ab_2+a_2)^2(aa_1b_0p+a_0a_1p-a_1cp-b_0)=0 \end{aligned} \quad (25)$$

$$\begin{aligned} & (ab_0+a_0-c)^2(2ab_2^2p+2ab_0b_2+ab_1^2+2a_2b_2p+2a_0b_2+a_1b_1-2b_2c)+2(ab_0+a_0-c)(ab_1+a_1)(3ab_1b_2p+ab_0b_1+2a_1b_2p+a_2b_1p+a_0b_1-b_1c+a_2)+(2(ab_0+a_0-c)(ab_2+a_2)+(ab_1+a_1)^2)(2ab_0b_2p+ab_1^2p+2a_0b_2p+a_1b_1p-2b_2cp+a_1)+2(ab_1+a_1)(ab_2+a_2)(ab_0b_1p+a_0b_1p-b_1cp+a_0)-ag(2ab_2+2a_2)=0 \end{aligned} \quad (26)$$

$$\begin{aligned} & (ab_0+a_0-c)^2(2aa_2b_2p+aa_1b_1+2aa_2b_0+2a_2^2p+2a_0a_2+a_1^2-2a_2c)+2(ab_0+a_0-c)(ab_1+a_1)(aa_1b_2p+2aa_2b_1p+aa_1b_0+3a_1a_2p+a_0a_1-a_1c-b_2)+(2(ab_0+a_0-c)(ab_2+a_2)+(ab_1+a_1)^2)(aa_1b_1p+2aa_2b_0p+2a_0a_2p+a_1^2p-2a_2cp-b_1)+2(ab_1+a_1)(ab_2+a_2)(aa_1b_0p+a_0a_1p-a_1cp-b_0)-g(2ab_2+2a_2)=0 \end{aligned} \quad (27)$$

利用 Maple 软件解非线性代数方程组(8)~(27)得到两组解:

$$\text{解组 1: } \{a = a, a_0 = a_0, a_1 = -ab_1, a_2 = -ab_2, b_0 = b_0, b_1 = b_1, b_2 = b_2, c = ab_0 + a_0\}$$

$$\text{解组 2: } \{a = a, a_0 = -ab_0 + c, a_1 = a_1, a_2 = 0, b_0 = b_0, b_1 = -\frac{a_1}{a}, b_2 = 0, c = c\}$$

将 2 个解组分别代入式(6)和式(2), 同时注意到方程(7)有解

$$\varphi(\xi) = \sqrt{p} \tan(\sqrt{p}\xi) = -\sqrt{p} \cot(\sqrt{p}\xi) \quad p > 0 \quad (28)$$

$$\varphi(\xi) = -\sqrt{-p} \tanh(\sqrt{-p}\xi) = \sqrt{-p} \coth(\sqrt{-p}\xi) \quad p < 0 \quad (29)$$

可得到方程(1)的精确行波解.

情形 1 将解组 1 代入式(6)和式(2), 得到洋流运动方程(1)有如下三角函数周期行波解和双曲函数行波解。

当  $p > 0$  时,

$$\begin{cases} U(\xi) = a_0 - ab_1 \sqrt{p} \tan(\sqrt{p}\xi) - ab_2 p \tan^2(\sqrt{p}\xi) \\ V(\xi) = b_0 + b_1 \sqrt{p} \tan(\sqrt{p}\xi) + b_2 p \tan^2(\sqrt{p}\xi) \\ H(\xi) = b_0 \xi + \frac{1}{2} b_1 \ln(1 + (\tan(\sqrt{p}\xi))^2) - b_2 p \xi + b_2 \sqrt{p} \tan(\sqrt{p}\xi) \end{cases} \quad (30)$$

当  $p < 0$  时,

$$\begin{cases} U(\xi) = a_0 + ab_1 \sqrt{-p} \tanh(\sqrt{-p}\xi) + ab_2 p (\tanh(\sqrt{-p}\xi))^2 \\ V(\xi) = b_0 - b_1 \sqrt{-p} \tanh(\sqrt{-p}) - b_2 p (\tanh(\sqrt{-p}))^2 \\ H(\xi) = b_0 \xi + \frac{1}{2} b_1 \ln(\tanh^2(\sqrt{-p}\xi) - 1) + \frac{b_2 p \tanh(\sqrt{-p}\xi)}{\sqrt{-p}} + \frac{b_2 p}{2\sqrt{-p}} \ln \frac{\tanh(\sqrt{-p}\xi - 1)}{\tanh(\sqrt{-p}\xi + 1)} \end{cases} \quad (31)$$

其中,  $\xi = x + ay - ct$ , 且  $c = ab_0 + a_0$ 。

情形2 将解组2代入式(6)和式(2), 得到洋流运动方程有(1)有如下三角函数周期行波解和双曲函数行波解。

当  $p > 0$  时,

$$\begin{cases} U(\xi) = (-ab_0 + c) - ab_1 \sqrt{p} \tan(\sqrt{p}\xi) \\ V(\xi) = b_0 + b_1 \sqrt{p} \tan(\sqrt{p}\xi) \quad p > 0 \\ H(\xi) = b_0 \xi + \frac{1}{2} b_1 \ln(1 + (\tan(\sqrt{p}\xi))^2) \end{cases} \quad (32)$$

当  $p < 0$  时,

$$\begin{cases} u(\xi) = (-ab_0 + c) - ab_1 \sqrt{-p} \tanh(\sqrt{-p}\xi) \\ V(\xi) = b_0 + b_1 \sqrt{-p} \tanh(\sqrt{-p}\xi) \quad p > 0 \\ H(\xi) = b_0 \xi - \frac{1}{2} b_1 \ln(\tanh^2(\sqrt{-p}\xi) - 1) \end{cases} \quad (33)$$

其中  $\xi = x + ay - ct$ 。

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## Exact Traveling Wave Solutions of the Oceanic Currents Motion Equations

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**Abstract:** We consider the oceanic currents motion equations. By using the generalized tanh-method and the symbolic toolbox of Maple, some explicit solutions of the oceanic currents motion equations are obtained, including the trigonometric functions periodic travelling wave solutions and the hyperbolic function travelling wave solutions.

**Keywords:** the oceanic currents motion equations; nonlinear PDEs; periodic wave solutions; hyperbolic function solutions